

# Invariant Polynomial Functions on $k$ qudits

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## Abstract

We study the polynomial functions on tensor states in  $(\mathbb{C}^n)^{\otimes k}$  which are invariant under  $SU(n)^k$ . We describe the space of invariant polynomials in terms of symmetric group representations. For  $k$  even, the smallest degree for invariant polynomials is  $n$  and in degree  $n$  we find a natural generalization of the determinant. For  $n, d$  fixed, we describe the asymptotic behavior of the dimension of the space of invariants as  $k \rightarrow \infty$ . We study in detail the space of homogeneous degree 4 invariant polynomial functions on  $(\mathbb{C}^2)^{\otimes k}$ .

## 1 Introduction

In quantum mechanics, a combination of states in Hilbert spaces  $H_1, \dots, H_k$  leads to a state in the tensor product Hilbert space  $H_1 \otimes \dots \otimes H_k$ . Such a state will be called here a tensor state. In this paper we take  $H_1 = \dots = H_k = \mathbb{C}^n$  where  $n > 1$ . Then a tensor state is a joint state of  $k$  qudits. It would be very interesting to classify tensor states in  $(\mathbb{C}^n)^{\otimes k}$  up to the action of the product  $U(n)^k$  of unitary groups of local symmetries. A natural approach to this is to study the algebra of invariant polynomials. This approach was developed by Rains [R], by Grassl, Rötteler and Beth [G-R-B1] [G-R-B2], by Linden and Popescu [L-P] and by Coffman, Kundu and Wootters [C-K-W]. These authors study the ring of invariant polynomials in the components of a tensor state in  $(\mathbb{C}^n)^{\otimes k}$  and in their complex-conjugates. For  $k$  qubits, explicit descriptions of invariants are given in [G-R-B1], [G-R-B2], [L-P] and in [C-K-W].

In this paper the symmetry group we consider is the product  $G = SU(n)^k$  of special unitary groups; one thinks of  $G$  as the special group of local symmetries. We study the  $G$ -invariant polynomial functions  $Q$  on the tensor states in  $(\mathbb{C}^n)^{\otimes k}$  (we discuss in §2 how this is relevant to the description of the  $G$ -orbits). We consider polynomials in the entries of a tensor state, in other words, holomorphic polynomials.

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Let  $\mathcal{R}_{n,k,d}$  be the space of homogeneous degree  $d$  polynomial functions on tensor states in  $(\mathbb{C}^n)^{\otimes k}$ . Let  $\mathcal{R}_{n,k,d}^G$  be the space of  $G$ -invariants in  $\mathcal{R}_{n,k,d}$ . See §2 for more discussion. We reduce the problem of computing  $\mathcal{R}_{n,k,d}^G$  to a problem in the invariant theory of the symmetric group  $\mathfrak{S}_d$  (Proposition 2.1). In particular,  $\mathcal{R}_{n,k,d}^G$  is non-zero only if  $d$  is a multiple of  $n$ . So the “first” case is  $d = n$ ; we examine this in §3. We find that if  $k$  is odd then  $\mathcal{R}_{n,k,n}^G = 0$  while if  $k$  is even then  $\mathcal{R}_{n,k,n}^G$  is 1-dimensional. In the latter case we write down (§3) explicitly the corresponding invariant polynomial  $P_{n,k}$  in  $\mathcal{R}_{n,k,n}$ ; we find  $P_{n,k}$  is a natural generalization of the determinant of a square matrix.

For fixed  $n, d$  the direct sum  $\oplus_k \mathcal{R}_{n,k,d}$  is an associative algebra. We study the asymptotic behavior of  $\dim \mathcal{R}_{n,k,d}^G$  as  $k \rightarrow \infty$  in §4. In §5, we specialize to the case of  $k$ -qubits, i.e.  $n = 2$ . We compute the dimension of the space  $\mathcal{R}_{2,k,4}^G$  of degree 4 invariants as well as the dimension of the space of invariants in  $\mathcal{R}_{2,k,4}^G$  under the natural action of  $\mathfrak{S}_k$ . We show that  $\oplus_k \mathcal{R}_{2,k,4}^{\mathfrak{S}_k \ltimes G}$  is a polynomial algebra on 2 generators. For  $k \leq 5$  we describe the representation of  $\mathfrak{S}_k$  on  $\mathcal{R}_{2,k,4}^G$ . For  $k = 4$  we find some interesting relations with the results on classification of tensor states in  $(\mathbb{C}^2)^{\otimes 4}$  given in [B].

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## 2 Polynomial invariants of tensor states

We will consider  $(\mathbb{C}^n)^{\otimes k}$  as a space of contravariant tensor states  $u$ . Then (once we fix a basis of  $\mathbb{C}^n$ )  $u$  is given by  $n^k$  components  $u^{p_1 p_2 \dots p_k}$ . We consider the algebra  $\mathcal{R}_{n,k}$  of polynomial functions on  $(\mathbb{C}^n)^{\otimes k}$ . So  $\mathcal{R}_{n,k}$  is the polynomial algebra  $\mathbb{C}[x_{p_1 p_2 \dots p_k}]$  in the  $n^k$  coordinate functions  $x_{p_1 p_2 \dots p_k}$ . We have a natural algebra grading  $\mathcal{R}_{n,k} = \oplus_{d=0}^{\infty} \mathcal{R}_{n,k,d}$  where  $\mathcal{R}_{n,k,d}$  is the space of homogeneous degree  $d$  polynomial functions.

A function in  $\mathcal{R}_{n,k,d}$  amounts to a symmetric degree  $d$  covariant tensor  $Q$  in  $(\mathbb{C}^n)^{\otimes k}$ . So  $Q$  has  $n^{dk}$  components  $Q_{i_{11} \dots i_{dk}}$  where we think of the indices  $i_{ab}$  as being arranged in a rectangular array of  $d$  rows and  $k$  columns and  $Q_{i_{11} \dots i_{dk}}$  is invariant under permutations of the rows of the array. Then  $Q$  defines the function

$$u \mapsto Q_{i_{11} \dots i_{dk}} u^{i_{11} i_{12} \dots i_{1k}} u^{i_{21} i_{22} \dots i_{2k}} \dots u^{i_{d1} i_{d2} \dots i_{dk}} \quad (2.1)$$

where we used the usual Einstein summation convention. In this way,  $\mathcal{R}_{n,k}$  identifies with  $S^d((\mathbb{C}^n)^{\otimes k})$ .

Now the group  $G = SU(n)^k$  acts on our tensor states  $u$  and tensors  $Q$  as follows. Let the matrix  $g_{ij}$  live in the  $m$ -th copy of  $SU(n)$  and let  $g^{ij}$  be the inverse matrix. Then  $g_{ij}$  transforms  $u^{p_1 p_2 \dots p_k}$  into  $g_{p_m q_m} u^{q_1 q_2 \dots q_k}$  and  $Q_{i_{11} \dots i_{dk}}$  into  $Q_{j_{11} \dots j_{dk}} g^{j_{1m} i_{1m}} g^{j_{2m} i_{2m}} \dots g^{j_{dm} i_{dm}}$ . The identification of  $\mathcal{R}_{n,k,d}$  with  $S^d((\mathbb{C}^n)^{\otimes k})$  is  $G$ -equivariant.

We are interested in the algebra  $\mathcal{R}_{n,k}^G = \oplus_{d=0}^{\infty} \mathcal{R}_{n,k,d}^G$  of  $G$ -invariants. We view this as a first step towards studying the orbits of  $G$  on  $(\mathbb{C}^n)^{\otimes k}$ . One can first study the orbits of the complex group  $G_{\mathbb{C}} = SL(n, \mathbb{C})^k$  and then decompose the  $G_{\mathbb{C}}$ -orbits under the  $G$ -action. Note that a polynomial is  $G$ -invariant if and only if it is  $G_{\mathbb{C}}$ -invariant. The closed  $G_{\mathbb{C}}$  orbits play a special role – they are the most degenerate orbits. Given any orbit  $Y$ , its closure contains a unique closed orbit  $Z$ ; then points in  $Y$  degenerate to

points in  $Z$ . The  $G_{\mathbb{C}}$ -invariant functions separate the closed orbits; they take the same values on  $Y$  and on  $Z$ . The set of closed orbits of  $G_{\mathbb{C}}$  in  $(\mathbb{C}^n)^{\otimes k}$  has the structure of an affine complex algebraic variety with  $\mathcal{R}_{n,k}^G$  as its algebra of regular functions. Thus a complete description of  $\mathcal{R}_{n,k}^G$  would lead to a precise knowledge of the closed  $G_{\mathbb{C}}$ -orbits.

Our approach is thus somewhat different from that of [R] [G-R-B1] [G-R-B2] [L-P] [C-K-W] who study the invariant functions on  $(\mathbb{C}^n)^{\otimes k}$  which are polynomials in the  $x_{p_1 \dots p_k}$  and in their complex conjugates; these can also be described as the invariant polynomial functions on  $(\mathbb{C}^n)^{\otimes k} \oplus \overline{(\mathbb{C}^n)^{\otimes k}}$ .

At this point it is useful to examine the case  $k = 2$ . We can identify  $(\mathbb{C}^n)^{\otimes 2}$  with the space  $M_n(\mathbb{C})$  of square matrices and then  $G = SU(n)^2$  acts on  $M_n(\mathbb{C})$  by  $(g, h) \cdot u = guh^{-1}$ . So  $\mathcal{R}_{n,k,d}^G$  is the space of homogeneous degree  $d$  polynomial functions  $Q$  of an  $n$  by  $n$  matrix  $u$  which are bi- $SL(n, \mathbb{C})$ -invariant, i.e.  $Q(guh^{-1}) = Q(u)$  for  $g, h \in SL(n, \mathbb{C})$ . Then  $Q$  is, up to scaling, the  $r$ th power of the determinant  $D$  for some  $r$ . Hence  $d = rn$ . It follows that  $\mathcal{R}_{n,2}^G$  is the polynomial algebra  $\mathbb{C}[D]$ . Thus the space of closed orbits for  $SL(n, \mathbb{C})^2$  identifies with  $\mathbb{C}$ , where  $\lambda$  corresponds to the unique closed orbit  $Z_\lambda$  inside the set  $X_\lambda$  of matrices of determinant  $\lambda$ . For  $\lambda \neq 0$ ,  $Z_\lambda = X_\lambda$  while for  $\lambda = 0$ ,  $Z_0$  reduces to the zero matrix.

We view  $S^d((\mathbb{C}^n)^{\otimes k})$  as the space of invariants for the symmetric group  $\mathfrak{S}_d$  acting on  $((\mathbb{C}^n)^{\otimes k})^{\otimes d}$ . So

$$\mathcal{R}_{n,k,d}^G = (((\mathbb{C}^n)^{\otimes k})^{\otimes d})^{G \times \mathfrak{S}_d} = (((\mathbb{C}^n)^{\otimes d})^{\otimes k})^{G \times \mathfrak{S}_d} \quad (2.2)$$

Recall the Schur decomposition  $(\mathbb{C}^n)^{\otimes d} = \bigoplus_{\alpha} S^{\alpha}(\mathbb{C}^n) \otimes E_{\alpha}$  where  $\alpha$  ranges over partitions of  $d$  with at most  $n$  rows,  $S^{\alpha}(\mathbb{C}^n)$  is the irreducible covariant representation of  $SU(n)$  given by the Schur functor  $S^{\alpha}$ , and  $E_{\alpha}$  is the corresponding irreducible representation of  $\mathfrak{S}_d$ . We use the convention that  $E_{\alpha}$  is the trivial representation if  $\alpha = [d]$ , while  $E_{\alpha}$  is the sign representation if  $\alpha = [1^d]$ . Thus we have

$$((\mathbb{C}^n)^{\otimes d})^{\otimes k} = \sum_{|\alpha_1| = \dots = |\alpha_k| = d} S^{\alpha_1}(\mathbb{C}^n) \otimes \dots \otimes S^{\alpha_k}(\mathbb{C}^n) \otimes E_{\alpha_1} \otimes \dots \otimes E_{\alpha_k} \quad (2.3)$$

Now taking the invariants under  $G \times \mathfrak{S}_d$  we get

$$\mathcal{R}_{n,k,d}^G = \sum_{|\alpha_1| = \dots = |\alpha_k| = d} S^{\alpha_1}(\mathbb{C}^n)^{SU(n)} \otimes \dots \otimes S^{\alpha_k}(\mathbb{C}^n)^{SU(n)} \otimes (E_{\alpha_1} \otimes \dots \otimes E_{\alpha_k})^{\mathfrak{S}_d} \quad (2.4)$$

The representation  $S^{\alpha_j}(\mathbb{C}^n)$ , since it is irreducible, has no  $SU(n)$ -invariants except if  $S^{\alpha_j}(\mathbb{C}^n) = \mathbb{C}$  is trivial. This happens if and only if  $\alpha_j$  is a rectangular partition with all columns of length  $n$ . This proves:

**Proposition 2.1.** *If  $n$  does not divide  $d$ , then  $\mathcal{R}_{n,k,d}^G = 0$ . If  $d = nr$ , then  $\mathcal{R}_{n,k,d}^G$  is isomorphic to  $(E_{\pi}^{\otimes k})^{\mathfrak{S}_d}$  where  $\pi = [r^n]$ .*

The permutation action of  $\mathfrak{S}_k$  on  $(\mathbb{C}^n)^{\otimes k}$  induces an action of  $\mathfrak{S}_k$  on  $\mathcal{R}_{n,k,d}^G$ .

**Corollary 2.2.** *The isomorphism of Proposition 2.1 intertwines the  $\mathfrak{S}_k$ -action on  $\mathcal{R}_{n,k,d}^G$  with the action of  $\mathfrak{S}_k$  on  $(E_{\pi}^{\otimes k})^{\mathfrak{S}_d}$  given by permuting the  $k$  factors  $E_{\pi}$ .*

### 3 The generalized determinant function

Given  $n$  and  $k$ , we want to find the smallest positive value of  $d$  such that  $\mathcal{R}_{n,k,d}^G \neq 0$ . By Proposition 2.1, the first candidate is  $d = n$ .

**Corollary 3.1.**  $\mathcal{R}_{n,k,n}^G \neq 0$  iff  $k$  is even. In that case,  $\mathcal{R}_{n,k,n}^G$  is one-dimensional and consists of the multiples of the function  $P_{n,k}$  given by

$$P_{n,k}(u) = \sum_{\sigma_2, \dots, \sigma_k \in \mathfrak{S}_n} \epsilon(\sigma_2) \cdots \epsilon(\sigma_k) \prod_{h=1}^n u^{h h_{\sigma_2} \cdots h_{\sigma_k}} \quad (3.1)$$

where  $h_{\sigma_j} = \sigma_j(h)$ .

*Proof.* By Proposition 2.1, we need to compute  $(E_\pi^{\otimes k})^{\mathfrak{S}_d}$ . For  $d = n$ ,  $\pi = [1^n]$  and so  $E_\pi$  is the sign representation of  $\mathfrak{S}_n$ . Then  $(E_\pi^{\otimes k})$  is one-dimensional and carries the trivial representation if  $k$  is even, or the sign representation if  $k$  is odd.

Now for  $k$  even, we can easily compute a non-zero function  $P = P_{n,k}$  in  $\mathcal{R}_{n,k,n}$ . For  $S^\pi(\mathbb{C}^n)$  is the top exterior power  $\wedge^n \mathbb{C}^n$ . Thus  $P$  is a non-zero element of the one-dimensional subspace  $(\wedge^n \mathbb{C}^n)^{\otimes k}$  of  $((\mathbb{C}^n)^{\otimes n})^{\otimes k}$ . The tensor components of  $P$  are then given by  $P_{i_{11} \cdots i_{nk}} = \frac{1}{n!} \epsilon(\sigma_1) \cdots \epsilon(\sigma_k)$  if for each  $j$ , the column  $i_{1j}, \dots, i_{nj}$  is a permutation  $\sigma_j$  of  $1, \dots, n$  and 0 otherwise. Then we get

$$P_{n,k}(u) = \frac{1}{n!} \sum_{\sigma_1, \dots, \sigma_k \in \mathfrak{S}_n} \epsilon(\sigma_1) \cdots \epsilon(\sigma_k) \prod_{h=1}^n u^{h_{\sigma_1} \cdots h_{\sigma_k}} \quad (3.2)$$

where  $h_{\sigma_i} = \sigma_i(h)$ . The expression is very redundant, as each term appears  $n!$  times. We remedy this by restricting the first permutation  $\sigma_1$  to be 1. This gives (3.1).  $\square$

$P_{n,k}$  is a *generalized determinant*;  $P_{n,k}$  is invariant under the  $\mathfrak{S}_k$ -action. For  $k = 2$ , (3.1) reduces to the usual formula for the matrix determinant.

Recall that the rank  $s$  of a tensor state  $u$  in  $(\mathbb{C}^n)^{\otimes k}$  is the smallest integer  $s$  such that  $u$  can be written as  $u = v_1 + v_2 + \cdots + v_s$ , where the  $v_i$  are decomposable tensor states  $v_i = w_{i1} \otimes w_{i2} \otimes \cdots \otimes w_{ik}$ . There is a relation between the rank and the vanishing of  $P_{n,k}$  as follows:

**Corollary 3.2.** If the tensor state  $u$  in  $(\mathbb{C}^n)^{\otimes k}$  has rank less than  $n$ , then  $P_{n,k}(u) = 0$ .

It is easy to find a tensor state  $u$  of rank  $n$  such that  $P_{n,k}(u)$  is non-zero. For instance,  $P_{n,k}(u) = 1$  if  $u$  has all components zero except  $u^{1 \cdots 1} = \cdots = u^{n \cdots n} = 1$ . For  $k = 2$ ,  $P_{n,k}(u) = 0$  implies  $u$  has rank less than  $n$ . For bigger (even)  $k$ , this is false, if  $n$  is large enough. This happens essentially because the rank of  $u$  can be very large (at least  $\frac{n^k}{kn-k+1}$ ). Thus  $P_{n,k}$  gives only partial information about the rank.

### 4 Asymptotics as $k \rightarrow \infty$

Suppose we fix  $n$  and  $d$  where  $d = rn$ . Then there is a  $G$ -invariant associative graded algebra structure  $P \circ Q$  on the direct sum  $\oplus_k \mathcal{R}_{n,k,d}^G$ . Indeed, the product of tensors

induces a  $(G \times \mathfrak{S}_d)$ -invariant map  $V^{\otimes k} \otimes V^{\otimes l} \rightarrow V^{\otimes(k+l)}$  where  $V = (\mathbb{C}^n)^{\otimes d}$ . The induced multiplication on the spaces of  $(G \times \mathfrak{S}_d)$ -invariants gives the product on  $\oplus_k \mathcal{R}_{n,k,d}^G$ , where we use the identification in (2.2). This multiplication corresponds, under the isomorphism of Proposition 2.1, to the product map  $E_\pi^{\otimes k} \otimes E_\pi^{\otimes l} \rightarrow E_\pi^{\otimes(k+l)}$ . This structure is very useful. For instance, if  $d = n$ , then  $P_{n,k} \circ P_{n,l} = \frac{1}{n!} P_{n,k+l}$ . Thus the determinant  $P_{n,2}$  determines  $P_{n,2m}$  in that the  $m$ -fold product  $P_{n,2} \circ \cdots \circ P_{n,2}$  is equal to  $(n!)^{-m+1} P_{n,2m}$ .

We will study the size of the algebra  $\oplus_k \mathcal{R}_{n,k,d}^G$  by finding an asymptotic formula for the dimension of  $\mathcal{R}_{n,k,d}^G$ . We do this for  $r \geq 2$ . Indeed for  $r = 1$  we already know  $\dim \mathcal{R}_{n,k,n}^G$  is 1 if  $k$  is even or 0 if  $k$  is odd; we call this the *static* case. The asymptotics involve the number

$$p = \dim E_\pi = d! \prod_{m=0}^{n-1} \frac{m!}{(m+r)!} \quad (4.1)$$

where  $\pi = [r^n]$  as in Proposition 2.1. Our formula for  $p$  is immediate from the hook formula for the dimension of an irreducible symmetric group representation.

**Proposition 4.1.** *Assume  $d = rn$  with  $r \geq 2$ . Then  $\dim \mathcal{R}_{n,k,d}^G \sim c \frac{p^k}{d!}$  as  $k \rightarrow \infty$ , where  $c = 1$  with one exception:  $c = 4$  if  $n = 2, d = 4$ .*

*Proof.* Let  $s = \dim \mathcal{R}_{n,k,d}^G = \dim (E_\pi^{\otimes k})^{\mathfrak{S}_d}$ . Then  $s = \frac{1}{d!} \sum_{\sigma \in \mathfrak{S}_d} \chi(\sigma)^k$  where  $\chi : \mathfrak{S}_d \rightarrow \mathbb{Z}$  is the character of  $E_\pi$ . If  $\sigma$  acts trivially on  $E_\pi$ , then  $\chi(\sigma) = p$ . If  $\sigma$  acts non-trivially, we claim  $|\chi(\sigma)| < p$ . To show this, it suffices to show that  $\sigma$  has at least two distinct eigenvalues on  $E_\pi$ ; this is because  $\chi(\sigma)$  is the sum of the  $p$  eigenvalues of  $\sigma$ . Now the set  $\mathfrak{T}_d$  of  $\sigma \in \mathfrak{S}_d$  which act on  $E_\pi$  by a scalar is a normal subgroup of  $\mathfrak{S}_d$ . So if  $d \geq 5$ , then  $\mathfrak{T}_d$  is  $\{1\}$ , the alternating group  $\mathfrak{A}_d$  or  $\mathfrak{S}_d$ . We can easily rule out the latter two possibilities, so  $\mathfrak{T}_d = \{1\}$ , which proves our claim. If  $d \leq 4$ , then (since  $r > 1$  and  $n > 1$ ), we have  $d = 4, n = 2$  and  $\pi = [2, 2]$ . Our claim is clear here since  $\mathfrak{S}_4$  acts on  $E_\pi$  through the reflection representation of  $\mathfrak{S}_3$  on  $\mathbb{C}^2$ .

Therefore we have  $s = c \frac{p^k}{d!} + o(p^k)$  as  $k \rightarrow \infty$  where  $c$  is cardinality of the kernel of  $\mathfrak{S}_d \rightarrow \text{Aut } E_\pi$ . Our work in the previous paragraph computes  $c$ .  $\square$

Proposition 4.1 implies that the algebra  $\oplus_k \mathcal{R}_{n,k,d}^G$  is far from commutative, as it has roughly  $1/N$  times the size of the tensor algebra  $\oplus_k (\mathbb{C}^p)^{\otimes k}$ . We note however that the  $\mathfrak{S}_k$ -invariants in  $\oplus_k \mathcal{R}_{n,k,d}^G$  form a commutative subalgebra, isomorphic to  $S(E_\pi)^{\mathfrak{S}_d}$ .

## 5 Quartic invariants of $k$ qubits

The case  $n = 2$  is of particular interest, as here the qudits are qubits, and this is the case being most discussed in quantum computation. Here we can give some precise non-asymptotic results for the first non-static case, namely  $\mathcal{R}_{2,k,4}^G$ . We put  $E = E_\pi = E_{[2,2]}$ . The proof of Proposition 4.1 easily gives

**Corollary 5.1.** *We have  $\dim \mathcal{R}_{2,k,4}^G = \frac{1}{3}(2^{k-1} + (-1)^k)$ .*

The first few values of  $\dim \mathcal{R}_{2,k,4}^G$ , starting at  $k = 1$ , are 0, 1, 1, 3, 5, 11, 21, 43. For  $k = 2$  and  $k = 3$  the unique (up to scalar) invariants are, respectively, the squared determinant  $P_{2,2}^2$  and the Cayley hyperdeterminant  $H_{2,3}$  (see [G-K-Z]). We note that the hyperdeterminant is very closely related to the relative tangle of 3 entangled qubits discussed in [C-K-W].

It would be useful to study  $\mathcal{R}_{2,k,4}^G$  as a representation of  $\mathfrak{S}_k$ , where  $\mathfrak{S}_k$  acts by permuting the  $k$  qubits. The  $\mathfrak{S}_k$ -invariants in  $\mathcal{R}_{2,k,4}^G$  are the  $(\mathfrak{S}_k \ltimes G)$ -invariants in  $\mathcal{R}_{2,k,4}$ . These  $(\mathfrak{S}_k \ltimes G)$ -invariant polynomials are very significant as they separate the closed orbits of the extended symmetry group  $\mathfrak{S}_k \ltimes SL(2, \mathbb{C})^k$  acting on  $(\mathbb{C}^2)^{\otimes k}$ . We can compute the dimension of the  $\mathfrak{S}_k$ -invariants as follows:

**Proposition 5.2.** *The dimension of the space of  $\mathfrak{S}_k \ltimes G$ -invariants in  $\mathcal{R}_{2,k,4}$  is  $M_k = \left\lfloor \frac{k}{6} \right\rfloor + r_k$  where  $r_k = 0$  if  $k \equiv 1 \pmod{6}$ , or  $r_k = 1$  otherwise. Furthermore the algebra  $\oplus_k \mathcal{R}_{2,k,4}^{\mathfrak{S}_k \ltimes G}$  is the polynomial algebra  $\mathbb{C}[P_{2,2}^2, H_{2,3}]$ .*

*Proof.* We have isomorphisms  $\mathcal{R}_{2,k,4}^{\mathfrak{S}_k \ltimes G} \simeq (E^{\otimes k})^{\mathfrak{S}_k \times \mathfrak{S}_3} \simeq S^k(E)^{\mathfrak{S}_3}$  since the representation of  $\mathfrak{S}_4$  on  $E$  factors through  $\mathfrak{S}_3$ . Thus the algebra  $\oplus_k \mathcal{R}_{2,k,4}^{\mathfrak{S}_k \ltimes G}$  identifies with  $S(E)^{\mathfrak{S}_3}$ . Now  $S(E)^{\mathfrak{S}_3}$  is the algebra of  $\mathfrak{S}_3$ -invariant polynomial functions on traceless  $3 \times 3$  diagonal matrices, and so is a polynomial algebra on the functions  $A \mapsto \text{Tr}(A^2)$  and  $A \mapsto \text{Tr}(A^3)$ . These invariants correspond (up to scaling) to  $P_{2,2}^2$  and  $H_{2,3}$ . The formula for the dimension follows easily.  $\square$

For instance, we have:  $M_1 = 0$ ,  $M_k = 1$  for  $2 \leq k \leq 5$ , and  $M_6 = 2$ . We remark that by replacing  $S(E)^{\mathfrak{S}_3}$  by  $\wedge(E)^{\mathfrak{S}_3}$ , it is easy to prove that the sign representation of  $\mathfrak{S}_k$  does not occur in  $(E^{\otimes k})^{\mathfrak{S}_4}$  for any  $k \geq 2$ .

We can determine the  $\mathfrak{S}_k$ -representation on  $\mathcal{R}_{2,k,4}^G$  for small  $k$  by explicit trace computations. For  $k = 2$  and  $k = 3$  we have the trivial 1-dimensional representation. For  $k = 4$ , we find  $\mathcal{R}_{2,4,4}^G$  is the direct sum  $E_{[4]} \oplus E_{[2,2]}$ . The trivial representation  $E_{[4]}$  of  $\mathfrak{S}_4$  is spanned by  $P_{2,4}^2$ , while the 2-dimensional representation  $E = E_{[2,2]}$  is spanned by the determinants  $\Delta(ijkl)$  introduced in [B]. Here  $(ijkl)$  is a permutation of  $(1234)$ . Given a tensor state  $u \in (\mathbb{C}^2)^{\otimes 4}$ , we can view it as an element  $v$  of  $\mathbb{C}^4 \otimes \mathbb{C}^4$ , where the first (resp. second)  $\mathbb{C}^4$  is the tensor product of the  $i$ -th and  $j$ -th copies of  $\mathbb{C}^2$  (resp. of the  $k$ -th and  $l$ -th copies). Then  $\Delta(ijkl)(u)$  is the determinant of  $v$ . As shown in [B], the  $\Delta(ijkl)$  span the representation  $E$  of  $\mathfrak{S}_4$ . The significance of the  $\Delta(ijkl)$  is that their vanishing describes the closure of the set of tensor states in  $(\mathbb{C}^2)^{\otimes 4}$  of rank  $\leq 3$ . For  $k = 5$  the representation  $\mathcal{R}_{2,5,4}^G$  of  $\mathfrak{S}_5$  is  $E_{[5]} \oplus E_{[2,1,1,1]}$ .

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